## ECE 456 - Problem Set 3 (Part 1)

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## Problem 1

(a) (i) The matrix equation is

$$[\hat{H}]\{\phi\} = E\{\phi\},\,$$

where  $[\hat{H}]$  is an N-by-N matrix with  $[\hat{H}]_{nm} = 0$  except for the following elements:

$$[\hat{H}]_{nn} = 2t_0 + U_n$$
$$[\hat{H}]_{n,n\pm 1} = -t_0$$
$$[\hat{H}]_{0,N} = [\hat{H}]_{N,0} = -t_0,$$

with  $t_0 = \hbar^2/(2ma^2)$  and  $U_n = U(na)$ . The N-vector  $\{\phi\}$  has elements  $\phi_n$  which each represent the value of the eigenvector at the point  $na = x_n$ .

(ii) The expression of the wave function  $\phi(x)$  as a sum of basis functions is as below:

$$\phi(x) = \sum_{n=1}^{N} \phi_n u_n(x)$$

The derived matrix equation:

$$[\hat{H}]_u\{\phi\} = [S]_u\{\phi\}$$

Where  $[\hat{H}]_u$  is a matrix with the elements:

$$H_{nm} = \int u_n^*(x) \hat{H} u_m(x) dx$$

and  $[S]_u$  is a matrix with the elements:

$$S_{nm} = \int u_n^*(x)u_m(x)dx$$

 $[\hat{H}]_u$  and  $[S]_u$  are both of size N-by-N. The elements of  $\{\phi\}$ ,  $\phi_n$ , are the expansion coefficients of  $\phi(x)$ .

**(b)** (i) code:

The bonding and antibonding eigenenergies are  $\boxed{-32.2567\,\mathrm{eV}}$  and  $\boxed{-15.5978\,\mathrm{eV}}$  respectively.

(ii) Neglecting normalization, we have the following expressions for  $\phi_B(z)$  and  $\phi_A(z)$ :

$$\phi_B(z) = u_L(z) + u_R(z)$$
  
$$\phi_A(z) = u_L(z) - u_R(z).$$

We obtain the following plot:

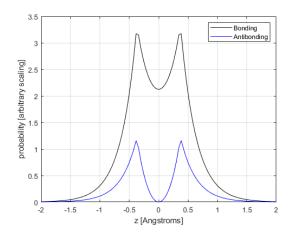


Figure 1. non-normalized probability densities for bonding and antibonding solutions.

## Problem 2

(a) (i)

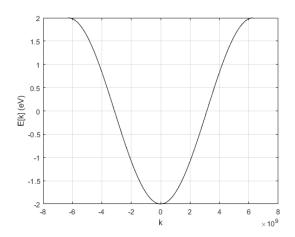


Figure 2. Energy vs. wave vector relationship (note: not discretized to account for N)

- (ii) Energy values from -2 eV to 2 eV are allowed.
- (iii) The vector  $\{\phi\}$ , which is of length N, and has elements n is:

$$\{\phi\} = Ce^{ik \cdot na}$$

The corresponding wave function is:

$$\phi(x) = \sum_{n=1}^{N} Ce^{ik \cdot na} u_n(x)$$

There is one wave function and thus one energy level for each value of k. This means that there is one electronic state per k.

**(b)** (i)

$$\{\phi\} = \begin{bmatrix} C_A e^{ika} \\ C_B e^{ika} \\ C_A e^{ik2a} \\ C_B e^{ik2a} \\ \vdots \\ C_A e^{ikNa} \\ C_A e^{ikNa} \end{bmatrix}$$

(ii)

$$\phi(x) = \sum_{n=1}^{N} C_A e^{ikna} u_{nA}(x) + C_B e^{ikna} u_{nB}(x)$$

(iii) [h(k)] is of size 2-by-2. Thus there will be two values of E(k) for a fixed k. This also means there are two  $\phi(x)$  for each k.

(c) (i)

$$\begin{split} \phi_2 + 2\phi_3 + \phi_4 &= E\phi_1 \\ \phi_1 + \phi_3 + 2\phi_4 &= E\phi_2 \\ 2\phi_1 + \phi_2 + \phi_4 &= E\phi_3 \\ \phi_1 + 2\phi_2 + \phi_3 &= E\phi_4 \end{split}$$

(ii)

$$\phi_2 + 2\phi_3 + \phi_0 = E\phi_1$$

$$\phi_1 + \phi_3 + 2\phi_4 = E\phi_2$$

$$2\phi_5 + \phi_2 + \phi_4 = E\phi_3$$

$$\phi_5 + 2\phi_6 + \phi_3 = E\phi_4$$

(iii) A generalized form of the *n*th equation is:

$$E\phi_n = \phi_{n-1} + \phi_{n+1} + 2\phi_{n+2}$$

(iv)

$$ECe^{ikna} = Ce^{ik(n+1)a} + Ce^{ik(n-1)a} + 2Ce^{ik(n+2)a}$$

$$E = e^{ika} + e^{-ika} + 2e^{2ika}$$

$$E(k) = 2e^{2ika} + 2\cos(ka)$$
(1)

(v) Imposing the repeating boundary conditions  $\phi_{n+4} = \phi_n$ , We obtain the following relationship:

$$Ce^{ikna} = Ce^{ik(n+4)a}$$
$$1 = e^{i4ka}$$

For this to hold, 4ka must be some multiple of  $2\pi$ , and this mean  $k = \frac{\pi}{2a} \cdot integer$ .

(vi) Using the E-k relationship from equation 1, we know that k must always be real, so the  $2\cos(ka)$  portion of the E-k relationship must be real. As well, if we substitute in the equation for k we obtained in part (v), we get the following (partial) expression:

$$2e^{i2ka} = 2e^{i\pi \cdot integer}$$

Which we know will always be real (with a value of  $\pm 2$ ).

(vii) Since  $e^{2\pi n} = 1$ , we have:

$$\phi_n(k + \frac{2\pi}{a}) = Ce^{i(k + \frac{2\pi}{a} \cdot nA)} = Ce^{i(k \cdot nA)}$$
$$\phi_n(k + \frac{2\pi}{a}) = \phi_n(k).$$

Therefore wavefunctions for which k is separated by  $\frac{2\pi}{a}$  are equivalent, and we only need consider the range  $k \in [-\frac{\pi}{a}, \frac{\pi}{a}]$ .